# The Crystal Structure of Cyclotrimethylene-trinitramine

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The structure of cyclotrimethylene-trinitramine (RDX),  $C_3H_6N_6O_6$ , has been refined from single-crystal neutron-diffraction data. The final weighted R index for 836 independent reflections is 0.021. The compound crystallizes in the orthorhombic space group *Pbca*, a = 13.182 (2), b = 11.574 (2), c = 10.709 (2) Å, Z = 8. The molecule consists of alternate CH<sub>2</sub> and N-NO<sub>2</sub> groups in a puckered ring. The environment of the carbon atoms is essentially tetrahedral, and the N-NO<sub>2</sub> groups are planar. The molecule possesses a plane of approximate mirror symmetry perpendicular to the plane defined by the three carbon atoms. The thermal motion may be described by rigid-body motion of the ring and separate rigid-body motion of the nitro groups.

### Introduction

The compound cyclotrimethylene-trinitramine,  $C_3H_6N_6O_6$ , also known as Cyclonite, or by the code designation RDX, is a well known explosive. There are two known polymorphic forms, designated RDX(I) and RDX(II) (McCrone, 1950). The structure of RDX(II) has not been determined because of the difficulty of obtaining and preserving well formed crystals, even for short periods, owing to its extreme instability. This paper describes the refinement of the structure of RDX(I); it will be referred to hereafter simply as RDX.

The crystal structure of RDX was first investigated by Hultgren (1936). He reported that it possesses orthorhombic symmetry, with eight molecules per unit cell and most probable space group  $V_h^{15}(Pbca)$ . A complete crystal morphology has been reported by Mc-Crone (1950). The crystal structure of RDX was reinvestigated by Harris & Reed (1959) using three-dimensional X-ray photographic data. They determined the complete crystal structure except for the hydrogen positions. The present study was undertaken to complete the crystal structure by determining the hydrogen positions and refining overall structural details.

## Experimental

The unit-cell dimensions were determined by a leastsquares fit to scattering angles measured by X-ray diffraction. For this, the scattering angles  $(2\theta)$  of the 30 most intense reflections having scattering angles in the range 35 to 60° were measured on an automated Picker 4-circle diffractometer with Mo K $\alpha$  radiation monochromated by a pyrolytic graphite plate. The unit-cell dimensions thus obtained are,  $a = 13 \cdot 182$  (2),  $b = 11 \cdot 574$ (2), c = 10.709 (2) Å, and  $D_x = 1.806$ . The observed density is reported as  $D_m = 1.816$  (Federoff & Sheffield, 1966) and the melting point is 205 °C (McCrone, 1950).

A crystal, approximately  $1.4 \times 1.2 \times 0.9$  mm, was used for this experiment. The crystal was mounted on a four-circle neutron diffractometer controlled by a computer (Alperin & Prince, 1970). The crystal was oriented with its a axis nearly, but not exactly, parallel to the  $\varphi$  axis, in order to avoid systematic effects due to multiple diffraction. The diffraction intensities were measured by the procedure described by Prince (1972). With a neutron wavelength of 1.232 Å and a limiting  $2\theta$  angle of 100° there are 1590 independent reflections. The intensity at the peak position of each reflection was compared with the intensities at background positions on either side, and if the peak intensity  $(I_p)$  did not exceed background  $(I_b)$  by  $2\sigma$ , where  $\sigma$  is given by  $\sigma = (I_p + I_b)^{1/2}$ , the reflection was considered to be unobserved. 836 of the 1590 independent reflections had observable intensities. The intensities of all observed reflections were converted to relative structure factors and each was assigned a standard deviation,  $\sigma_F$ , based on counting statistics. The calculated linear absorption coefficient is 1.2 cm<sup>-1</sup>, assuming an effective incoherent scattering cross-section for hydrogen of  $42 \times$ 10<sup>-24</sup> cm<sup>2</sup> (Rush, Taylor & Havens, 1962). The maximum variation in F due to absorption is less than  $\pm 0.8$ %. This was considered to be negligible, and no absorption corrections were made.

The positional parameters for the heavy atoms, as determined by Harris & Reed (1959), were used to compute signs for  $F_o$  and  $(F_o - F_c)$  Fourier syntheses. All hydrogen atom positions were determined from the difference synthesis. Three cycles of least-squares refinement with the program *RFINE* (Finger, 1968) with isotropic temperature factors, gave a weighted *R* index  $\{wR = [\sum w(s|F_o| - |F_c|)^2 / \sum wF_c^2]^{1/2}$  where *s* is the scale factor and  $w = 1/\sigma_F^2\}$  of 0.093. The refinement was continued

with anisotropic temperature factors for all atoms. An isotropic secondary extinction parameter, S, was included in the refinement, using the formula (Zachariasen, 1968)  $F_{cor} = |F_c|(1 + S\beta F_c^2)^{-1/4}$ , where  $\beta$  is a parameter proportional to the Lorentz factor. The final value of S was  $0.484 \times 10^{-3}$  for  $\beta = 1/\sin 2\theta$ . The two strongest reflections, 102 and 200, with calculated extinction factors of 0.7282 and 0.7275 respectively, were omitted from the final stage of refinement because of severe extinction. The final discrepancy indices were wR =0.021 and R = 0.039 for all observed reflections except the two strongest reflections. The largest parameter shift in the last cycle was less than 0.01 of the standard deviation. The final least-squares parameters are given in Table 1, and the observed and calculated structure factors are listed in Table 2. There was no observable trend in the values of the function  $1/s\sum_{s} w(|F_{o}| - |F_{c}|)^{2}$ , where s is the number of reflections in a subset, for subsets of the observed reflections divided according to ranges of F and  $\sin \theta / \lambda$ . The difference Fourier map obtained after final least-squares refinement was featureless.

#### Discussion

The heavy atom positions obtained in the present work all agree well with those reported in the previous X-ray study of Harris & Reed (1959), with differences of less than 0.02 Å in bond length and 0.8° in bond angle. The interatomic distances and angles in the RDX molecule, as obtained in the present study, are given in Table 3. The molecule of RDX consists of three N-NO<sub>2</sub> groups linked to three methylene groups alternately to form a puckered, six-membered C-N ring (see Figs. 1 and 2). The N(1)-NO<sub>2</sub> group is essentially

Table 1. Final least-squares parameters for the structure of RDXThermal parameters are the coefficients,  $B_{ij}$  (Å<sup>2</sup>),

in the form exp  $\left[-\frac{1}{4} \left(B_{11}a^{*2}h^2 + B_{22}b^{*2}k^2 + B_{33}c^{*2}l^2 + 2B_{12}a^*b^*hk + 2B_{13}a^*c^*hl + 2B_{23}b^*c^*kl\right)\right]$ 

The estimated standard deviations are in parentheses.

	x/a ( × 104)	$y/b$ ( $\times 10^{4}$ )	z/c (×104)	$B_{11}$ ( $ imes 10^2$ )	$B_{22}$ (×10 <sup>2</sup> )	$B_{33}$ ( $\times 10^2$ )	$B_{12}~(\times 10^2)$	$B_{13}$ ( $\times 10^2$ )	$B_{23}$ ( $\times 10^2$ )
<b>C</b> (1)	1839 (3)	3578 (3)	4400 (3)	322 (15)	297 (13)	206 (14)	-70 (13)	-69 (11)	7 (12)
$\tilde{C}(2)$	503 (2)	2440 (2)	3395 (3)	233 (15)	145 (11)	344 (13)	8 (13)	-1(13)	-17 (11)
Č(3)	1487 (2)	3813 (3)	2159 (3)	252 (14)	332 (16)	175 (14)	-34(13)	34 (11)	-31 (12)
N(1)	1761 (2)	4360 (2)	3330 (2)	267 (8)	266 (8)	251 (9)	- 109 (8)	1 (8)	0 (8)
N(2)	877 (2)	2996 (2)	4537 (2)	295 (9)	210 (8)	250 (9)	14 (8)	33 (8)	9 (7)
N(3)	536 (1)	3222 (2)	2329 (2)	211 (9)	228 (8)	274 (9)	15 (8)	-35(7)	-37(7)
N(4)	2260 (2)	5379 (2)	3346 (2)	217 (8)	241 (9)	348 (10)	- 55 (8)	79 (9)	- 26 (9)
N(5)	155 (2)	3525 (2)	5294 (2)	455 (14)	322 (10)	256 (10)	157 (10)	104 (9)	77 (9)
N(6)	-333(2)	3875 (2)	2078 (2)	284 (11)	306 (10)	316 (10)	72 (9)	-72 (9)	- 53 (8)
O(1)	2270 (3)	5932 (3)	2388 (4)	566 (23)	330 (17)	407 (19)	-122 (18)	82 (17)	68 (16)
O(2)	2649 (3)	5687 (3)	4338 (4)	342 (16)	330 (17)	417 (19)	-86 (15)	-13 (16)	-131 (14)
O(3)	- 693 (3)	3142 (4)	5262 (4)	333 (21)	718 (26)	426 (20)	138 (21)	153 (17)	156 (19)
O(4)	454 (4)	4271 (3)	5990 (3)	894 (32)	381 (21)	258 (16)	251 (22)	57 (20)	- 53 (15)
O(5)	-1121 (3)	3534 (4)	2504 (4)	246 (16)	572 (25)	584 (25)	44 (17)	-62 (16)	22 (21)
O(6)	-236(4)	4691 (4)	1390 (4)	644 (27)	441 (23)	415 (20)	185 (21)	-91 (20)	62 (18)
H(1)	2401 (6)	2944 (7)	4226 (7)	372 (33)	532 (36)	564 (37)	-14 (30)	-128 (31)	141 (32)
H(2)	2013 (6)	4068 (6)	5244 (6)	762 (43)	560 (37)	270 (29)	-201 (34)	- 127 (29)	- 59 (28)
H(3)	-261(5)	2117 (6)	3540 (6)	371 (34)	433 (33)	597 (35)	- 146 (29)	-18 (31)	16 (27)
H(4)	1015 (5)	1728 (5)	3192 (7)	464 (30)	263 (26)	594 (37)	89 (27)	-16 (27)	-114 (27)
H(5)	2052 (5)	3166 (6)	1917 (6)	314 (28)	602 (37)	524 (37)	46 (30)	80 (25)	-225 (30)
H(6)	1440 (6)	4432 (7)	1413 (7)	637 (44)	602 (39)	363 (35)	-170 (35)	17 (28)	108 (31)



Fig. 1. RDX molecule showing a stereoscopic pair viewed from top of the molecular ring.

coplanar but the N(2)-NO<sub>2</sub> and N(3)-NO<sub>2</sub> groups are slightly bent. The distances of atoms from the best fit mean plane of each N-NO<sub>2</sub> group are given in Table 4. The bond lengths and bond angles of the N(2)-

 $NO_2$  and  $N(3)-NO_2$  groups are very similar but the configuration of the  $N(1)-NO_2$  group is quite different. For the  $N(1)-NO_2$  group, the N-N distance is shorter by approximately 0.04 Å, and the C-N-N angles are

# Table 2. Observed and calculated structure factors for RDX Unobserved reflections are indicated by an asterisk.

										70 12 82*		
H+2+1 0 466 471	10 72• -20 11 75• 3	0 373 373 1 130 -134 2 65= -45 3 82 -78	H,3,3 1 398 -392 2 191 -187 3 57* 39	Hr13r3 1 92 -65 2 73• -7	3 351 -350 4 69* -15 5 458 -461 6 209 217	0 660 -18 1 112 -101 2 94 102 3 169 164	6 206 213 7 202 -201 8 74* 11 9 74* 14	8 114 -124 9 142 -145 10 77• 71 11 227 222	9 286 -297 9 286 -297 10 860 -135 11 760 -82	8 82* -88 9 195 -174 10 78* 23 11 83* -123	в 73• -29 9 80• 37 н.5.10 1 155 165	H+5,12 1 790 -11 2 163 -149 3 380 15

larger by more than  $3^{\circ}$ . The mean plane of the N(1)-NO<sub>2</sub> group is inclined to the plane defined by the three carbon atoms by only 18°, but the  $N(2)-NO_2$  and N(3)-NO<sub>2</sub> planes are both inclined by 62°. Thus the trigonal molecular symmetry, which might be expected in the free RDX molecule, is not observed in the crystalline state, but there is a pseudo mirror plane passing through the atoms N(4), N(1), C(2), H(3), and H(4). These atoms are essentially coplanar with the plane defined by 0.1869x - 0.1101y + 0.9762z = 3.3648. The atomic distances from the pseudo mirror plane and the inclination angles of the interatomic vectors between the two pseudo mirror symmetric atoms are summarized in Table 5. Only oxygen atoms show significant departures from the mirror symmetric positions. All the pseudosymmetry-related bond lengths and angles in a molecule agree within two standard deviations, with the exception of the N-N-C angles, which disagree by at most  $1.5^{\circ}$ .

Table	3.	Interatomic	distances	and	angles	in	RDX
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(a) Distances			
C(1)-N(1) = 1	·464 (4) Å	N(1)-N(4)	1·351 (3) Å
C(1)-N(2) = 1	·443 (4)	N(2) - N(5)	1.392 (3)
C(2)-N(3) = 1	·468 (4)	N(3)–N(6)	1.398 (3)
C(2)-N(3) 1	•458 (4)		
C(3)-N(3) 1	·440 (4)	N(4)–O(1)	1.209 (5)
C(3)-N(1) 1	·450 (4)	N(4)-O(2)	1.233 (4)
		N(5)–O(3)	1.203 (5)
C(1)-H(1) = 1	·058 (10)	N(5)–O(4)	1.207 (5)
C(1)-H(2) 1	·092 (8)	N(6)–O(5)	1.201 (5)
C(2)-H(3) 1	·085 (8)	N(6)–O(6)	1.205 (5)
C(2)-H(4) 1	·087 (7)		
C(3)-H(5) = 1	·088 (8)		
C(3)-H(6) 1	·075 (9)		
(b) Angles			
N(1) - C(1) - N(2)	$107.8(2)^{\circ}$	C(3) - N(1) - C(1)	115·1 (2)°
N(1)-C(1)-H(1)	109.9 (4)	C(3) N(1) - N(4)	120.9 (2)
N(1)-C(1)-H(2)	110.0 (4)	C(1) - N(1) - N(4)	119.7 (2)
N(2)-C(1)-H(1)	108.0 (4)	N(1) - N(4) - O(1)	117.2 (3)
N(2)-C(1)-H(2)	110.0 (5)	N(1)-N(4)-O(2)	117.8 (3)
H(1)-C(1)-H(2)	111.0 (6)	O(1) N(4)-O(2)	125.0 (3)
N(2)-C(2)-N(3)	111.7 (2)	C(1)-N(2)-C(2)	114.6 (2)
N(2)-C(2)-H(3)	110.1 (4)	C(1)-N(2)-N(5)	117.1 (2)
N(2)-C(2)-H(4)	106.9 (4)	C(2)-N(2)-N(5)	116.6 (2)
N(3)-C(2)-H(3)	110.7 (4)	N(2)-N(5)-O(3)	117-2 (3)
N(3)-C(2)-H(4)	107·2 (4)	N(2)-N(5)-O(4)	116.8 (3)
H(3)-C(2)-H(4)	110-1 (6)	O(3) - N(5) - O(4)	125.7 (4)
N(3)-C(3)-N(1)	108.4 (2)	C(2)-N(3)-C(3)	114.8 (2)
N(3)-C(3)-H(5)	107.4 (4)	C(2)-N(3)-N(6)	117.5 (2)
N(3)-C(3)-H(6)	111.1 (4)	C(3) - N(3) - N(6)	115.6 (2)
N(1)-C(3)-H(5)	109.6 (4)	N(3)-N(6)-O(5)	117.3 (3)
N(1)-C(3)-H(6)	111.3 (5)	N(3)-N(6)-O(6)	117.0 (3)
H(5)-C(3)-H(6)	108.8 (6)	O(5) - N(6) - O(6)	125.5 (4)

The methylene groups have essentially tetrahedral configurations, with the bond angles in the range of  $109.5 \pm 2^{\circ}$ , and the H–C–H planes perpendicular to the N–C–N planes. The C–H bonds range from 1.06 to 1.09 Å. When corrections for the thermal vibrations are applied by the riding model, they increase to between 1.09 and 1.13 Å.

# Table 4. The atomic distances from the best-fit mean planes of N-NO<sub>2</sub> groups

The mean planes are:

0.8473x	-0.4601y	-0.2654z = -1.2993	for	$N(1)-NO_2$
0.2116x	-0.6696y	+0.7120z = 1.3744	for	$N(2)-NO_2$
-0.1646x	-0.5725v	-0.80327 = -4.2606	for	$N(3) - NO_2$

	N(1)-NO <sub>2</sub>	N(2)-NO <sub>2</sub>	N(3)-NO <sub>2</sub>	
N(1)	-0.002			
N(4)	0.008			
O(1)	-0.003			
O(2)	-0.003			
N(2)		0.007		
N(5)		-0.026		
O(3)		0.010		
O(4)		0.009		
N(3)			0.006	
N(6)			-0.022	
O(5)			0.008	
O(6)			0.008	
C(1)	0.198	-0.580		
C(2)		-0.536	-0.386	
C(3)	0.316		-0.446	

Table 5. Atomic distances from the pseudo mirror plane

Atoms in plane		Symmetry-related atom pair						
	Distances		Dista	nces	Angle			
N(4)	—0·005 Å	$C(1) \cdots C(3)$	−1·23 Å	1∙23 Å	90°			
N(1)	0.006	$N(2) \cdot \cdot \cdot N(3)$	-1.21	1.21	90			
C(2)	0.003	$N(5) \cdots N(6)$	- 1.76	1.77	89			
H(3)	-0.002	$O(2) \cdots O(1)$	-1.10	1.07	87			
H(4)	-0.005	$O(3) \cdots O(5)$	- 1.56	1.47	88			
		$O(4) \cdots O(6)$	- <b>2</b> ·46	2.57	89			
		$H(1) \cdots H(5)$	-1·27	1.26	90			
		$H(2) \cdot \cdot \cdot H(6)$	2.09	2.10	90			

One hydrogen atom in each methylene group is always located very close to oxygen atoms of adjacent nitro groups on both sides, with the distances ranging from 2.22 to 2.30 Å. The C...O intramolecular distances are also short, ranging from 2.63 to 2.68 Å. Although the C-H...O angles,  $99 \pm 1^{\circ}$ , are much smaller than those in conventional hydrogen bonds, the oxygen atoms of the nitro-groups are bound tightly to the molecular ring by strong O...H as well as O...C inter-



Fig. 2. Side view of the RDX molecule.

actions and tend to form approximately planar  $C_2N$ - $NO_2$  groups. Consequently, we may expect to find some degree of rigidity in the RDX molecule. The rigid body motion analysis by the method of Schomaker & Trueblood (1968) was applied to all the heavy atoms of the molecule, as shown in Table 6. The r.m.s. DELTA U(I,J) was 0.008 Å<sup>2</sup> and the e.s.d. of the DELTA U(I,J) was 0.009. The disagreement with rigid-body molecular motion was greatest for the nitro group atoms. When the rigid-body analysis was applied to the ring atoms alone, the agreement was excellent; the r.m.s. DELTA U(I,J) was 0.0013 Å<sup>2</sup> and the e.s.d. of the DELTA U(I,J) was 0.0019. These facts may suggest that each nitro group possesses an independent libration superimposed upon the molecular libration, as is observed in the other nitro compounds, cyclotetramethylene tetranitramine (Choi & Boutin, 1970) and s-trinitrobenzene (Choi & Abel; 1972). The principal axis of the largest molecular libration lies almost exactly on the pseudo mirror plane, and has inclination angles of  $34^{\circ}$  with the N(1)–N(4) bond and 56° with both the other N-N bonds. Rigid body motion analysis of the nitro group alone is not possible because of the indeterminacy which occurs in the case of planar groups with five or fewer atoms. However, an investigation of the thermal ellipsoids of the nitro-group atoms shows clear evidence of rigid body motion, as shown in Table 7. We define Cartesian coordinate axes (radial, normal, and tangent) to which we refer the directions of the largest root-mean-square displacements,  $R_3$ , separately for each atom. For each atom of a nitro group, the radial axis is directed from the ring nitrogen to which the nitro group is attached toward that atom, and the normal axis is perpendicular to both the radial axis and the  $O \cdots O$  vector of the nitro group. The normal axis, as defined above, is approximately normal to the best fit mean plane of the N–NO<sub>2</sub> group, with angular deviations of less than  $1.5^{\circ}$ for all three groups. The largest thermal vibration amplitudes occur in the directions of tangential axes for all atoms except O(1) and O(5), where they are in the normal direction. The direction of the least vibration amplitude is always along the radial axis. This is consistent with rigid-body motion of the nitro group with two dominant modes, one about the normal axis passing through the ring nitrogen atom and the other about a radial axis somewhere near the N-N bond.

### Table 6. Rigid-body motion analysis

(a) Rigid-body motion applied for all heavy atoms $\mathbf{T} = \begin{pmatrix} 0.028 & 0.0009 & 0.0030 \\ 0.0284 & -0.0005 \\ 0.0226 \end{pmatrix} \mathbf{L} = \begin{pmatrix} 0.0031 & -0.0004 & 0.0010 \\ 0.0064 & 0.0004 \\ 0.0039 \end{pmatrix}$
Principal axes of L (°) and their direction cosines
$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$
Displacements of libration axes from the center of mass (Å)
$\begin{array}{ll} L_{1,2} = & 0.057 & L_{1,3} = -0.261 \\ L_{2,1} = & -0.424 & L_{2,3} = -0.490 \\ L_{3,1} = & -0.505 & L_{3,2} = -0.705 \end{array}$
(b) Rigid-body motion applied for ring atoms alone
$\mathbf{T} = \begin{pmatrix} 0.0216 & 0.0015 & 0.0020 \\ 0.0206 & -0.0042 \\ 0.0248 \end{pmatrix} \mathbf{L} = \begin{pmatrix} 0.0068 & 0.0013 & 0.0009 \\ 0.0086 & -0.0023 \\ 0.0083 \end{pmatrix}$

Principal axes of L (°) and their direction cosines

-				
	6.0	0.0988	0.7505	-0.6535
	5.1	0.7722	0.3563	0.5260
	4∙0	<b>0</b> ∙6 <b>2</b> 76	-0.5566	-0.5443

Displacements of libration axes from the center of mass (Å)

$L_{1,2} = -0.289$	$L_{1,3} = 0.163$
$L_{2,1} = -0.088$	$L_{2,3} = -0.512$
$L_{3,1} = 0.229$	$L_{3,2} = -0.025$

There are several short intermolecular distances in RDX. Those considered to be the most significant are shown in Fig. 3.

It may be interesting to compare the crystal structure of RDX with those of both  $\alpha$ - and  $\beta$ -HMX (cyclotetramethylene-tetranitramine) (Cady, Larson & Cromer, 1963; Choi & Boutin, 1970), in view of the close similarity in chemical composition and explosive nature. The bond lengths and angles of RDX agree well with those of  $\alpha$ - and  $\beta$ -HMX, except for the C-N-C angles, which are approximately 8° smaller in RDX than in HMX. In both compounds, the N-NO<sub>2</sub> group planes are bound tightly to the molecular ring by strong  $O \cdots H$  as well as  $O \cdots C$  interactions. The conformation of the N-NO<sub>2</sub> group attachment to the molecular ring in RDX is very similar to that of  $\alpha$ -HMX. Very short intermolecular distances occur in both compounds.

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Table 7. Principal axes of the thermal ellipsoids and orientations of the longest axes

	$R_1$	$R_2$	$R_3$	Radial	Normal	Tangent	В
N(4)	0·142 (4) Å	0·177 (3) Å	0·225 (3) Å	91°	73°	163°	2.69 Å <sup>2</sup>
O(1)	0.170 (6)	0.238 (5)	0.282(6)	95	31	120	4.34
O(2)	0.156 (7)	0.218 (6)	0·257 (5)	86	73	163	3.63
N(5)	0.161 (4)	0.167 (4)	0.277 (4)	96	87	7	3.44
O(3)	0.167 (6)	0.226(6)	0.329 (6)	96	100	12	4.93
<b>O</b> (4)	0.158 (7)	0.207 (6)	0.356 (6)	87	91	3	5.11
N(6)	0.164 (4)	0.181 (3)	0.234 (3)	92	86	176	3.02
O(5)	0.170 (6)	0.269 (6)	0.276(6)	76	164	82	4.67
O(6)	0.183 (7)	0.246 (6)	0.310 (6)	88	107	163	5.00



Fig. 3. A diagram of the RDX molecule and some atoms of neighboring molecules showing short intermolecular distances.

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